

Precision measurements of W- and Z-boson transverse momentum spectra at ATLAS

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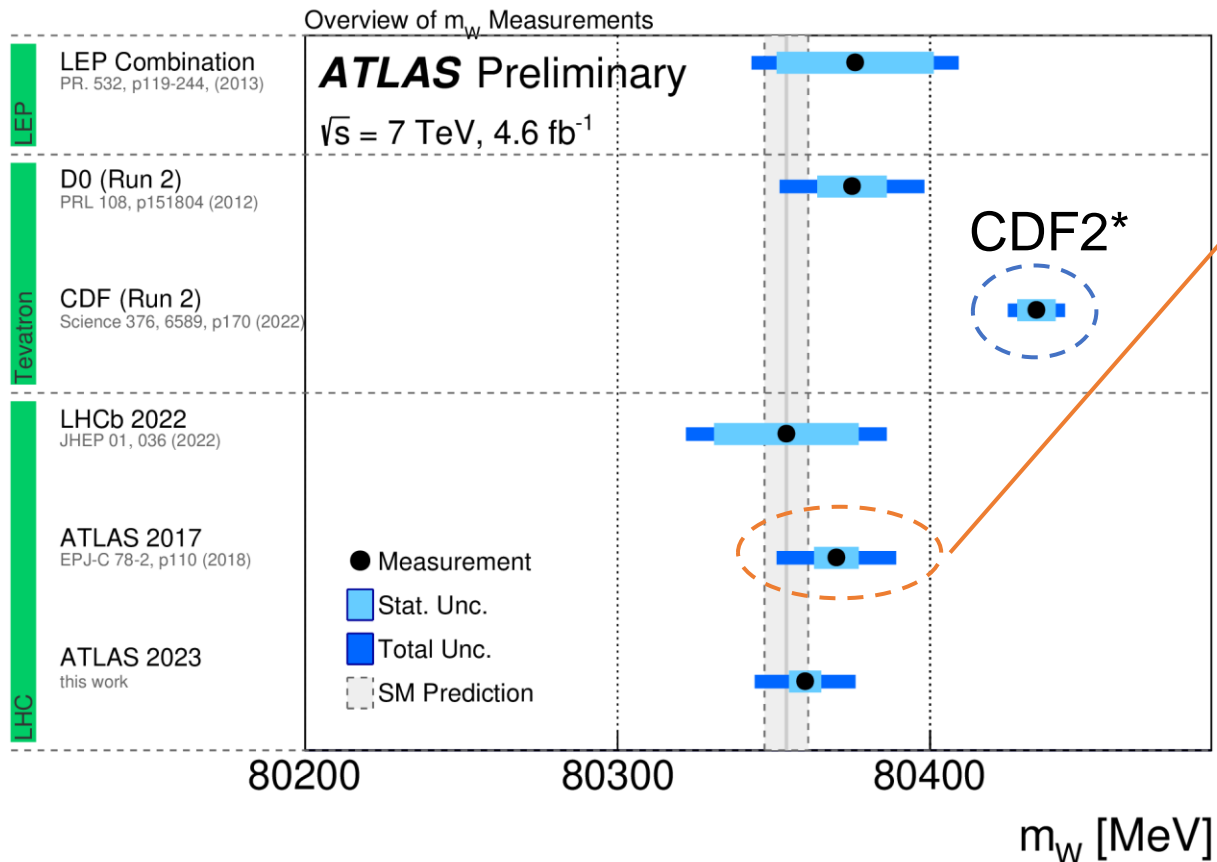
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The motivation of p_T^W and p_T^Z measurements

Reduce the p_T^W modelling uncertainty in the m_W measurement.



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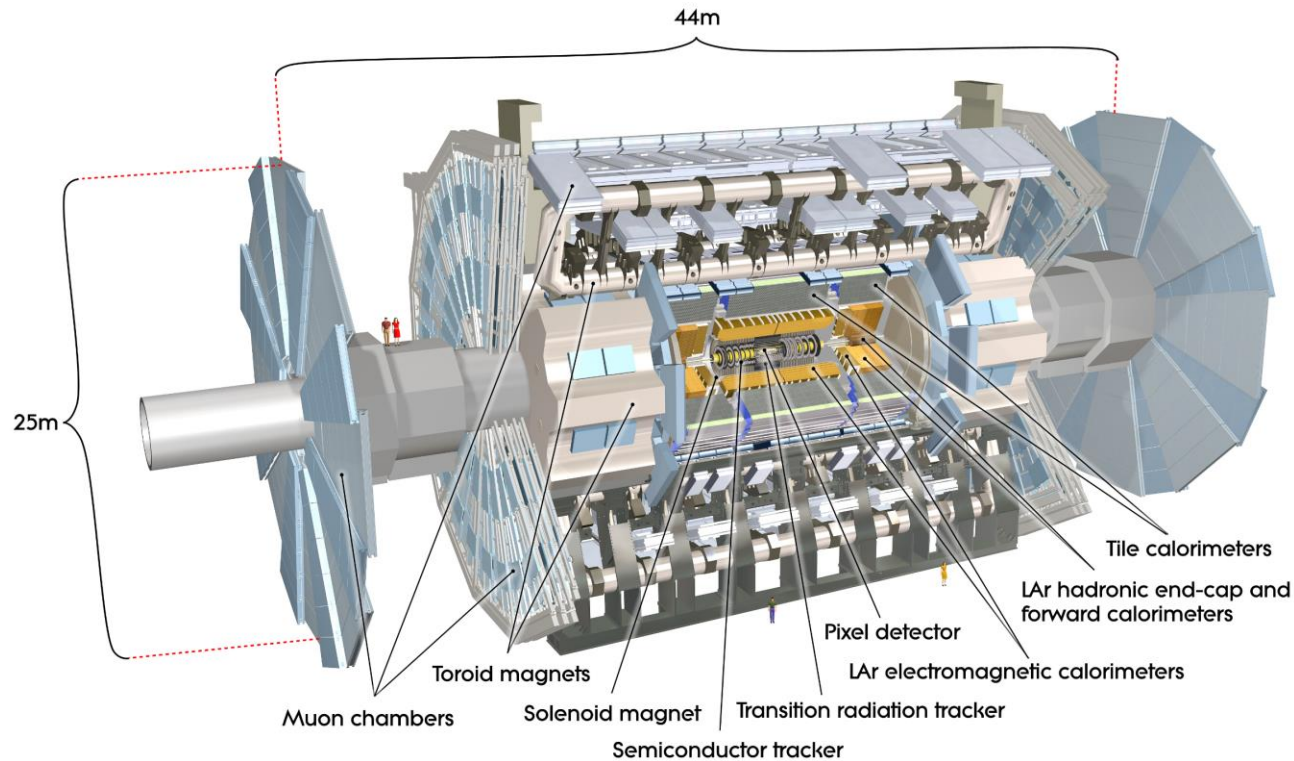
- ATLAS 7 TeV m_W analysis:
6 MeV p_T^W modelling uncertainty due to the $p_T^Z \rightarrow p_T^W$ extrapolation.
- Solution: **Direct p_T^W measurement.**
- Measuring both p_T^W and p_T^Z tests the underlying QCD predictions.

Targeted p_T^W precision:

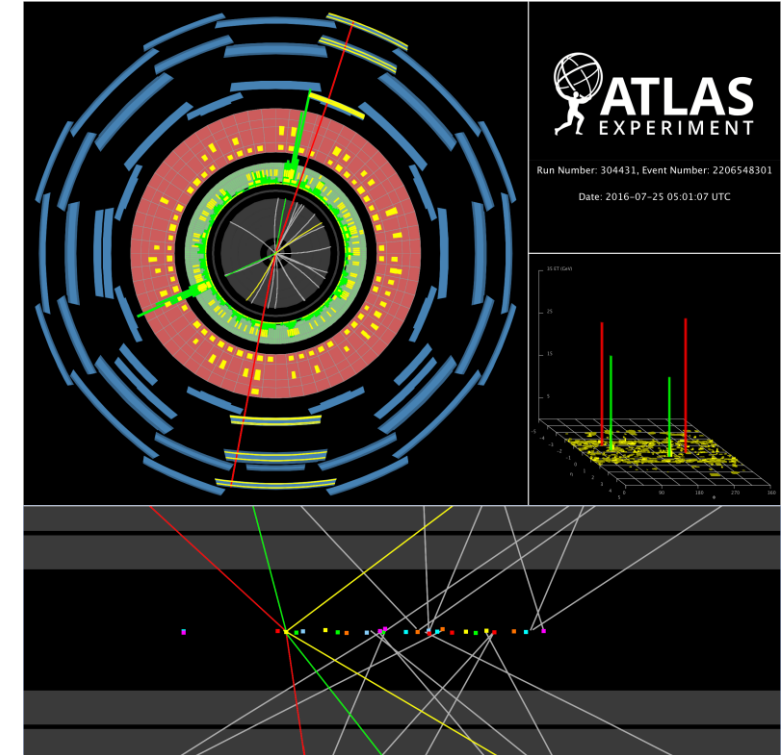
- A granularity of 6~7 GeV.
- 1~2% uncertainty for $p_T^W < 25 \text{ GeV}$ where the fixed-order perturbative prediction fails.

The ATLAS experiment

ATLAS detector: one of the general purpose detectors at LHC



Cut-away view of ATLAS



ATLAS detects the final state particles of proton-proton collisions.

e.g. $pp \rightarrow H \rightarrow ee\mu\mu$ event display

Event topology: leptonic decay of W/Z

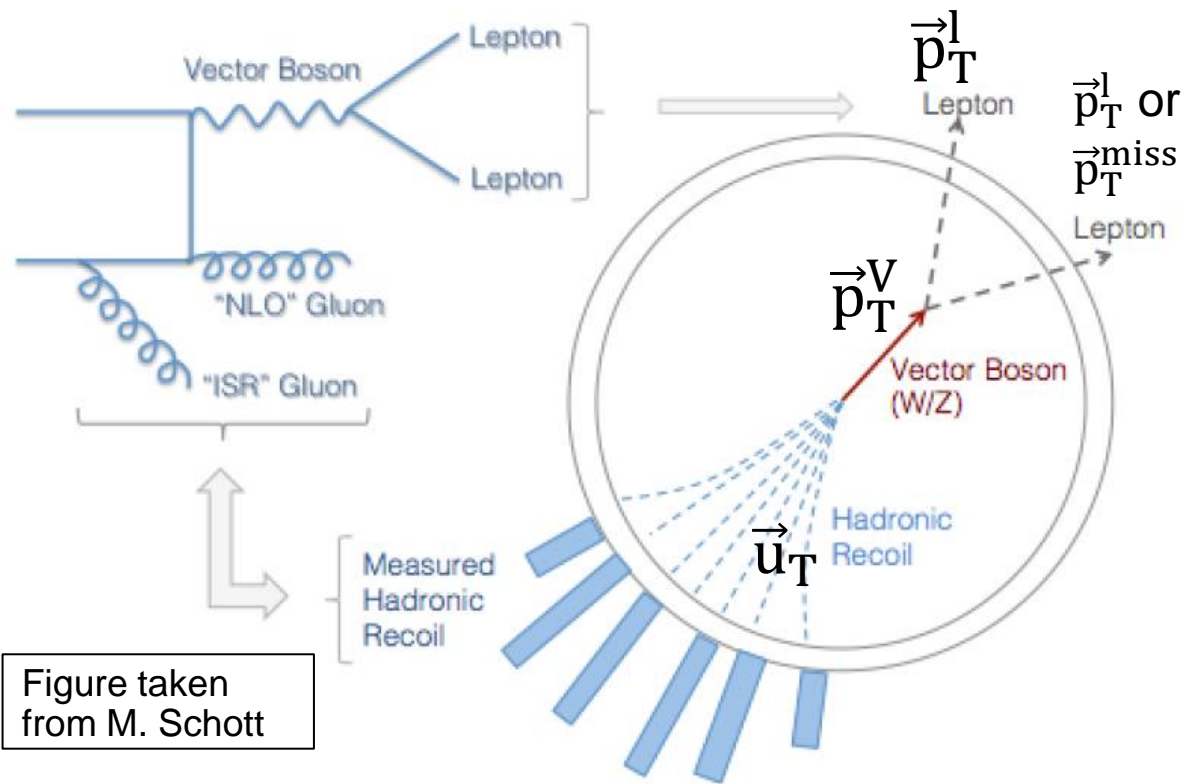


Figure taken from M. Schott

- Only **two objects** to measure:
 - (1) The charged lepton
 - (2) Hadronic recoil $\vec{u}_T = \sum \vec{p}_T^{\text{ISR } q,g} = -\vec{p}_T^V$

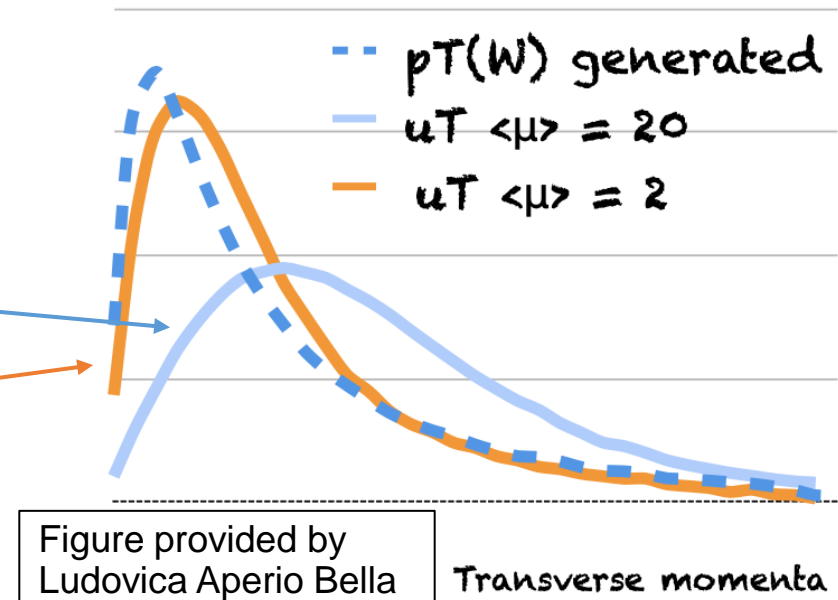
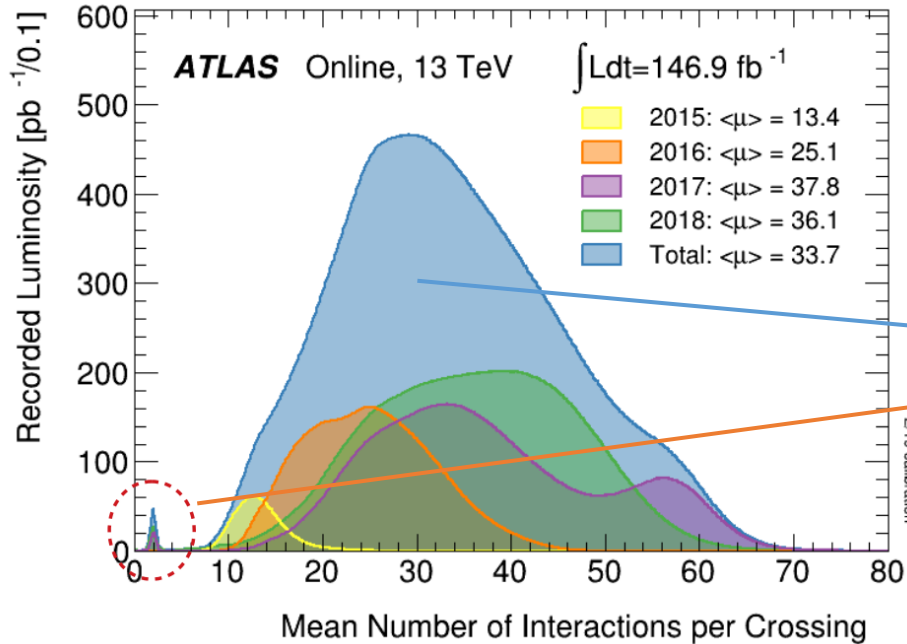
W events: $u_T \rightarrow p_T^W$

Z events: $p_T^{ll} \rightarrow p_T^Z$ or $u_T \rightarrow p_T^Z$

- Detector calibrations are needed for lepton and recoil.
- Detector resolution of u_T is affected by underlying event and pile-up.

**Lower pile-up in the dataset
 -> More precise measurement of p_T^W**

ATLAS Run 2 low pile-up data



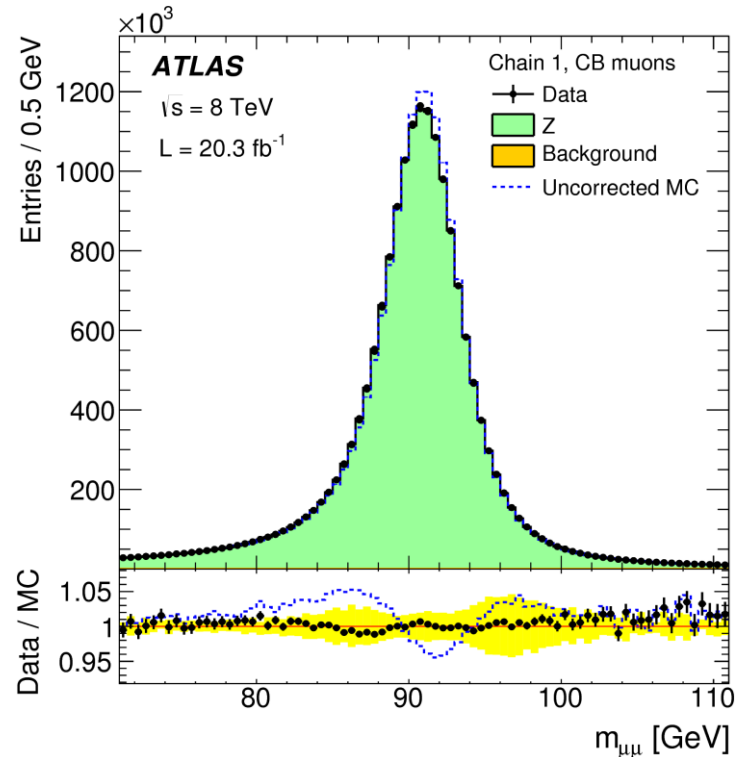
Probes of perturbative and non-perturbative QCD in W events, as well as in Z events at 5.02 TeV.

	2017, $\sqrt{s}=5.02 \text{ TeV}$	2017+2018, $\sqrt{s}=13 \text{ TeV}$
Luminosity (pb^{-1})	255	338
Total W events after selection	1.45M	4.36M
Total Z events after selection	122K	379K

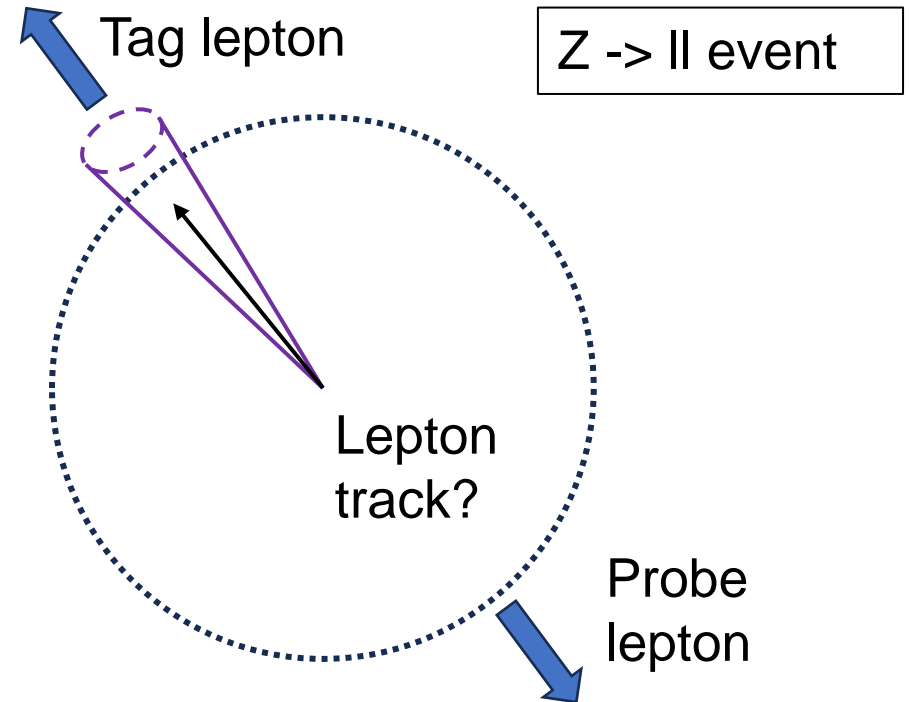
MC simulations matching the low pile-up condition in data:

- W & Z, top-related background and di-boson background.

Detector calibration

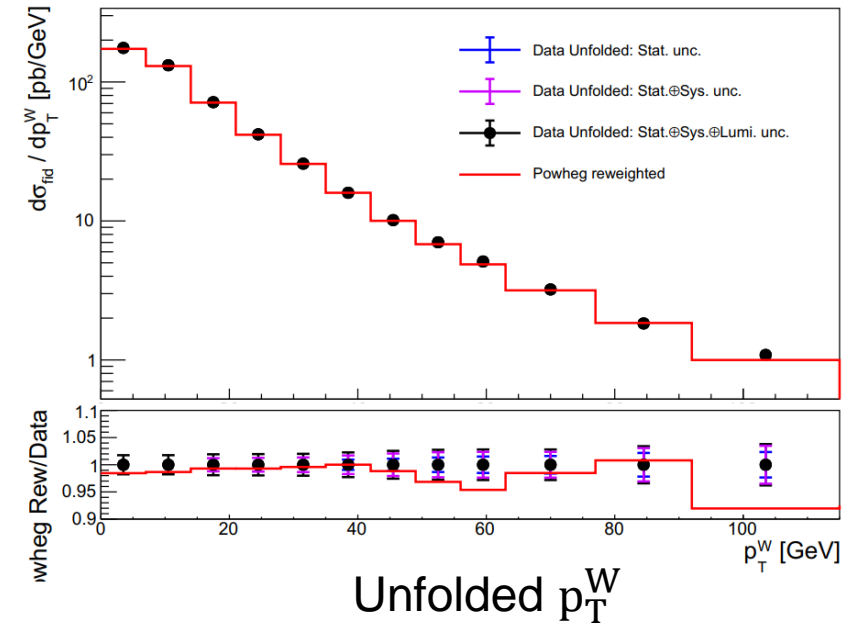
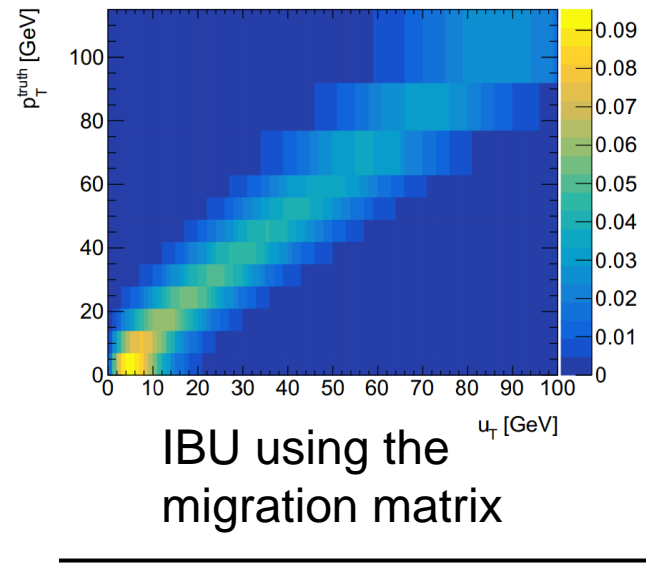
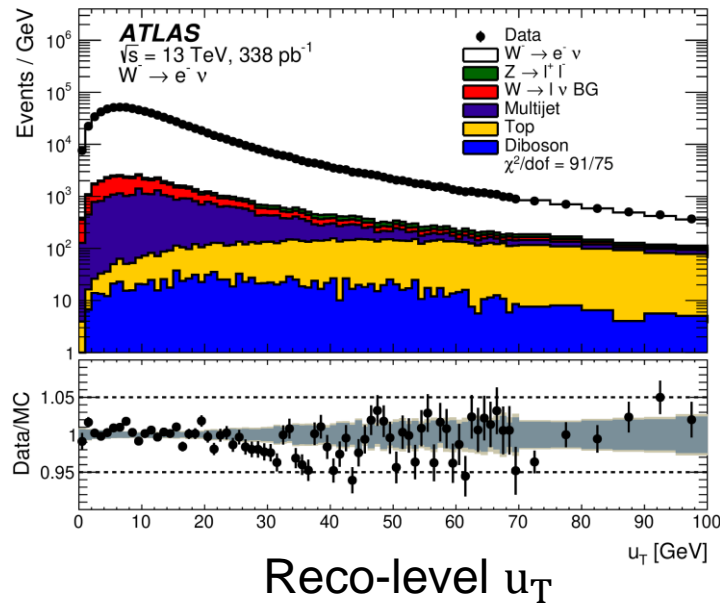


- The lepton momentum in the simulation is corrected to reproduce the resonance of Z-boson in data.
- Recoil calibration: use p_T^{ll} to constrain the detector response of recoil ($p_T^Z = p_T^{\text{ll}}$ and $\vec{u}_T + \vec{p}_T^V = 0$)



Efficiency measured in Z->ll events with “tag & probe”.

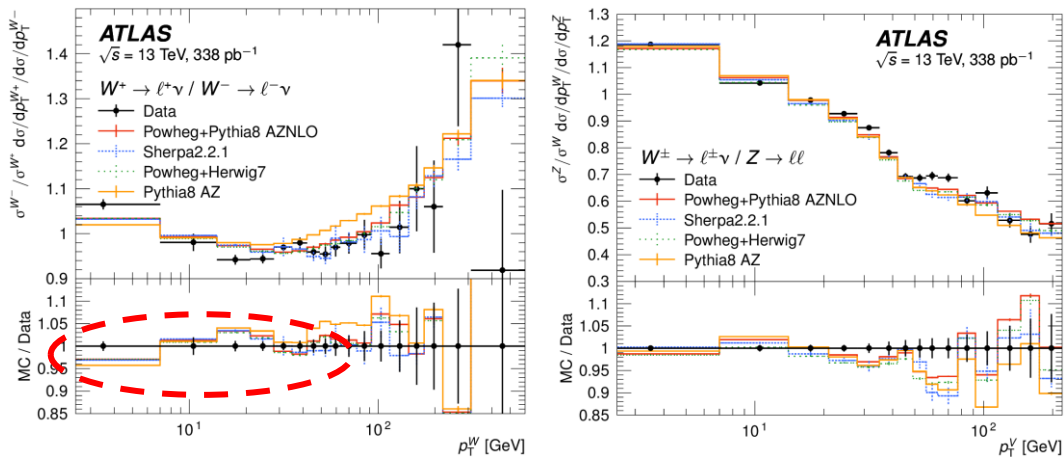
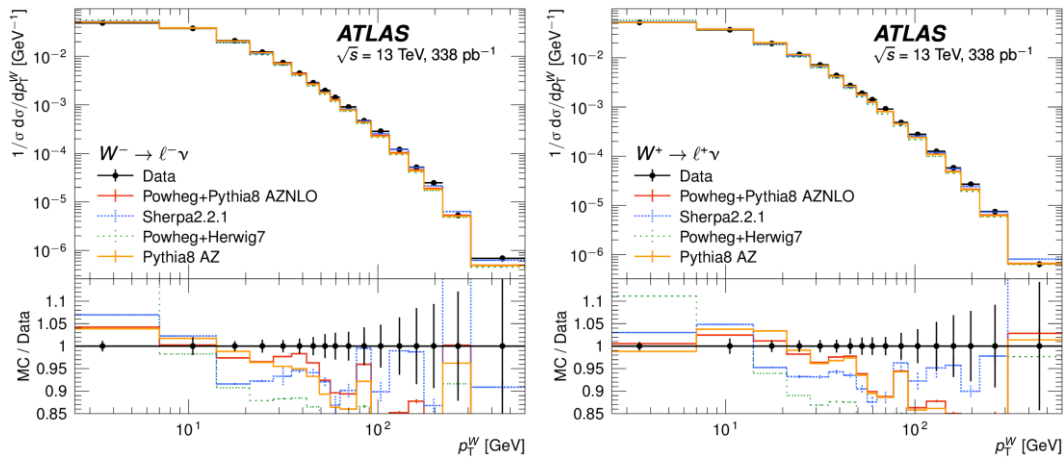
Unfolding



- Iterative Bayesian unfolding (IBU) for p_T^W : $\vec{u}_T = \Sigma \vec{p}_T^{\text{ISR } q,g} = -\vec{p}_T^V$
- p_T^Z spectrum can be obtained by unfolding u_T and p_T^{ll} .
- > Compatibility test of the unfolding

Results of the measurements

Differential cross-section @ 13 TeV



Integrated fiducial cross-section

Process	$\sigma_{\text{fid}}(\sqrt{s} = 5.02 \text{ TeV})$ [pb]	$\sigma_{\text{fid}}(\sqrt{s} = 13 \text{ TeV})$ [pb]
$W^- \rightarrow \ell^- \nu$	$1384 \pm 2 \text{ (stat.)} \pm 5 \text{ (syst.)} \pm 15 \text{ (lumi.)}$	$3486 \pm 3 \text{ (stat.)} \pm 18 \text{ (syst.)} \pm 34 \text{ (lumi.)}$
$W^+ \rightarrow \ell^+ \nu$	$2228 \pm 3 \text{ (stat.)} \pm 8 \text{ (syst.)} \pm 23 \text{ (lumi.)}$	$4571 \pm 3 \text{ (stat.)} \pm 21 \text{ (syst.)} \pm 44 \text{ (lumi.)}$
$Z \rightarrow \ell\ell$	$333.0 \pm 1.2 \text{ (stat.)} \pm 2.2 \text{ (syst.)} \pm 3.3 \text{ (lumi.)}$	$780.3 \pm 2.6 \text{ (stat.)} \pm 7.1 \text{ (syst.)} \pm 7.1 \text{ (lumi.)}$

Integrated cross-section ratios

Processes	Cross-section ratio at $\sqrt{s} = 5.02 \text{ TeV}$	Cross-section ratio at $\sqrt{s} = 13 \text{ TeV}$
W^+/W^-	$1.609 \pm 0.003 \text{ (stat.)} \pm 0.004 \text{ (syst.)}$	$1.308 \pm 0.003 \text{ (stat.)} \pm 0.004 \text{ (syst.)}$
W^-/Z	$4.16 \pm 0.02 \text{ (stat.)} \pm 0.03 \text{ (syst.)}$	$4.46 \pm 0.02 \text{ (stat.)} \pm 0.05 \text{ (syst.)}$
W^+/Z	$6.69 \pm 0.02 \text{ (stat.)} \pm 0.05 \text{ (syst.)}$	$5.84 \pm 0.02 \text{ (stat.)} \pm 0.06 \text{ (syst.)}$
W^\pm/Z	$10.85 \pm 0.04 \text{ (stat.)} \pm 0.07 \text{ (syst.)}$	$10.30 \pm 0.04 \text{ (stat.)} \pm 0.10 \text{ (syst.)}$

Paper submitted to EPJC ([arXiv:2404.06204](https://arxiv.org/abs/2404.06204))

The p_T^W uncertainties are propagated to the m_W fit for the preliminary result.

Summary

- The p_T^W measurement in ATLAS using special low pile-up data achieves remarkable precision (1-2% total uncertainty and a granularity of 7 GeV) and provides valuable input to the m_W measurement.
- The production of the Z-boson is compared with that of the W-boson at two well-separated center-of-mass energies.
- In addition to electro-weak precision measurements, I was also involved in the detector production for the upgrade of the ATLAS muon system* and the development of statistics method for high energy physics experiments**.
- Nowadays, I'm focusing on Higgs self-coupling and flavour tagging in ATLAS, as well as the detector simulation for FCC.

* [doi : 10.1016/j.nima.2021.166143](https://doi.org/10.1016/j.nima.2021.166143)

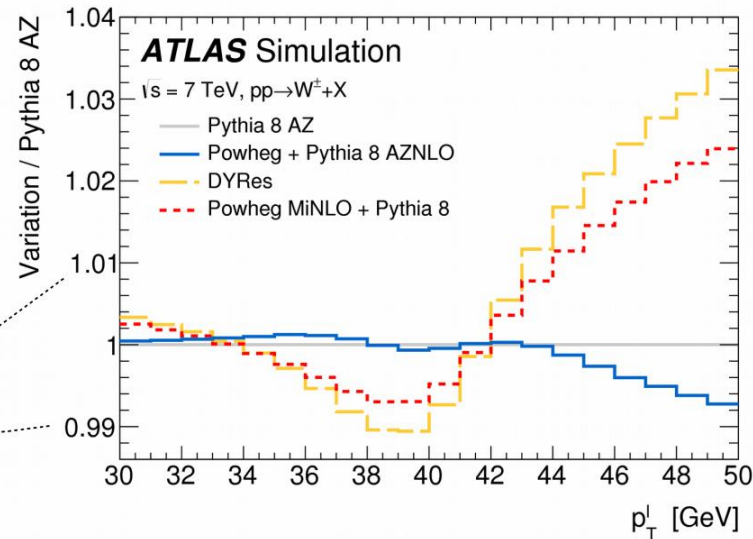
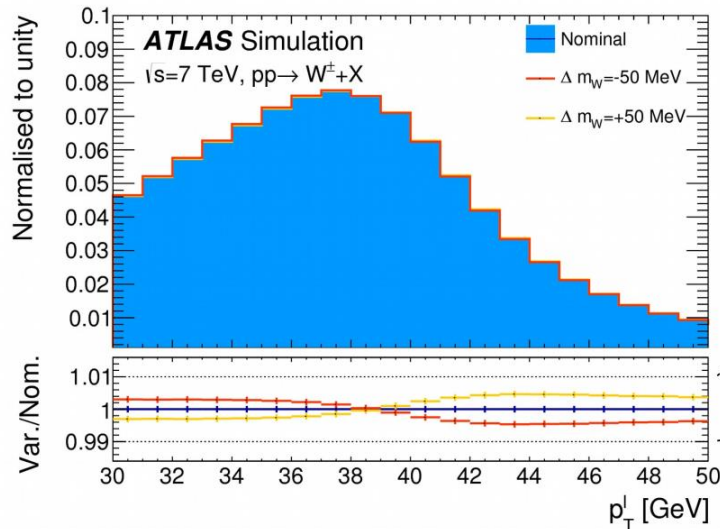
** [doi : 10.1140/epjc/s10052-024-12877-5](https://doi.org/10.1140/epjc/s10052-024-12877-5)

Thank you!

Backup

Motivation of the measurement

- In the m_W measurement, the lepton p_T spectrum requires a modelling of $p_T^W < \sim 1\%$ in the low values of p_T^W where the fixed-order perturbative prediction fails.
- Direct measurement of p_T^W , instead of modelling p_T^W based on measured p_T^Z , avoids the uncertainty due to the extrapolation.



- p_T^Z is also measured: Measuring both p_T^W and p_T^Z tests the differences in W and Z production processes. The measurement at 5.02 TeV is the input to the PDF fit.

Physics modelling

The “master formula” of differential cross-section for the Drell-Yan process

$$\frac{d\sigma}{dp_1 dp_2} = \left[\frac{d\sigma(m)}{dm} \right] \left[\frac{d\sigma(y)}{dy} \right] \left[\frac{d\sigma(p_T)}{dp_T} \Big|_y \frac{d\sigma(y)^{-1}}{dy} \right] \times \left[(1 + \cos^2\theta) + \sum_{i=0}^7 A_i(p_T, y, m) P_i(\cos\theta, \phi) \right]$$

For the leptonic decay of W/Z bosons, the kinematics of the decay product can be described by a factorization of differential cross-sections.

- $\frac{d\sigma(m)}{dm}$: The resonance peak of the vector boson -> Breit Wigner distribution.
- $\frac{d\sigma(y)}{dy}$: Boson rapidity
- $\frac{d\sigma(p_T)}{dp_T} \Big|_y \left(\frac{d\sigma(y)}{dy} \right)^{-1}$: Boson transverse momentum at given rapidity.
- A_i : Angular coefficients describe the polarization of the vector boson.

MC samples and event selection

In order to match the low pile-up condition in data:

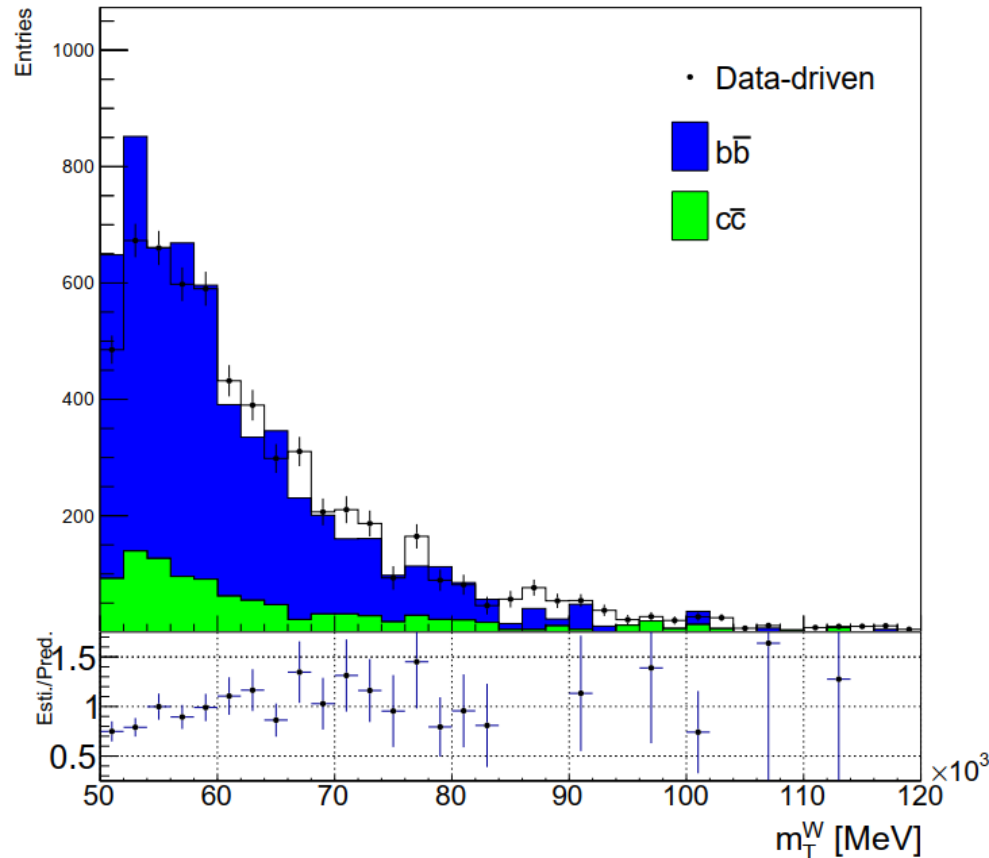
- W & Z , top-related background and di-boson background: MC simulations.
- Multijet background: Heavy flavor quark decays and pions faking electrons are hard to simulate accurately.
 - > Data-driven estimation of the multijet yield and shape.

Cut	Description
One charged lepton	Exactly one electron or muon
Lepton trigger	<ul style="list-style-type: none">• 1 electron, $E_T > 15$ GeV, loose ID.• Or 1 muon, $E_T > 14$ GeV.
Isolation	$Ptcone20 / \text{Min}(p_T^l, 50\text{GeV}) < 0.1$
Kinematics	$p_T^l > 25$ GeV
	$E_T^{\text{miss}} > 25$ GeV
	$m_T > 50$ GeV

Analysis cuts for W signal selection

ptconeXX: The sum of p_T of tracks in the given cone around the interested object.

Multijet background



The main contribution of multijet background in muon channels: $b\bar{b}$, $c\bar{c}$ decay. MC simulation normalized to the data-driven estimation.

In W events:

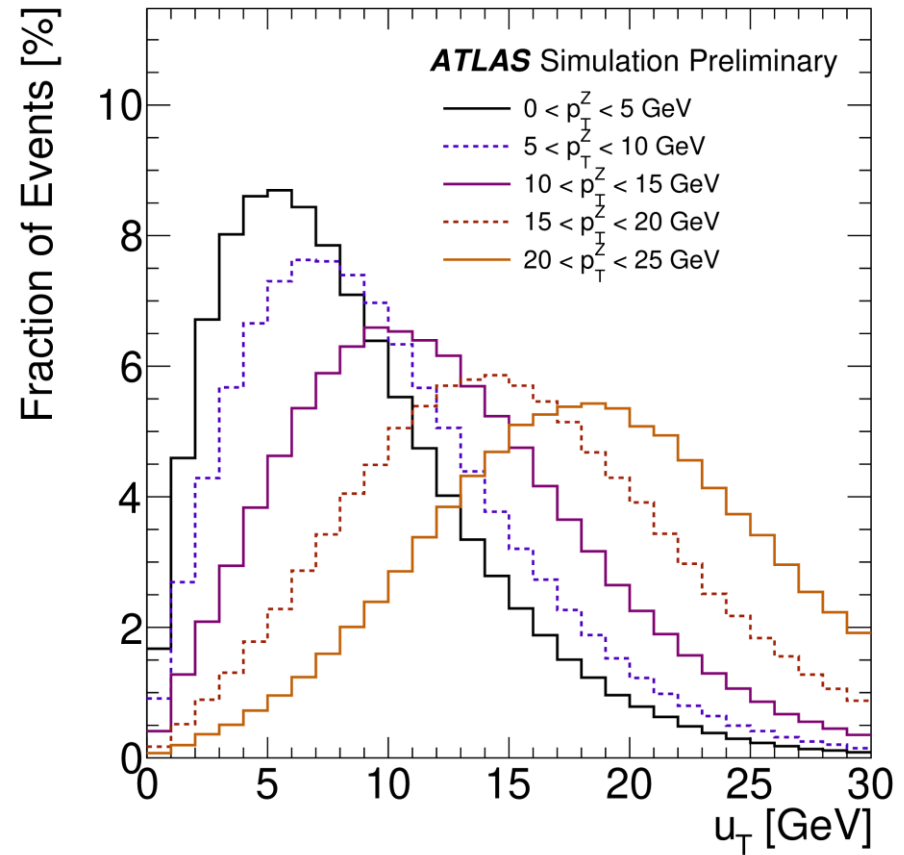
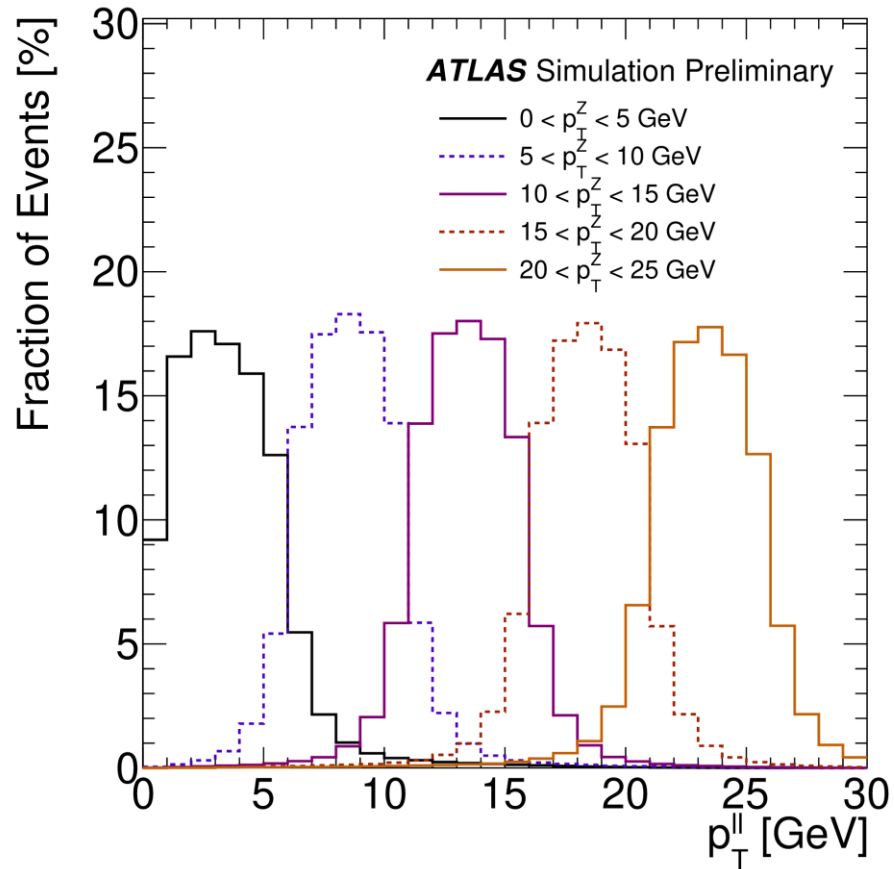
$$m_T = \sqrt{2p_T^l p_T^{\text{miss}} \left(1 - \cos(\phi_l - \phi_{p_T^{\text{miss}}})\right)}$$
$$\vec{p}_T^{\text{miss}} = -(\vec{p}_T^l + \vec{u}_T) \text{ for the neutrino}$$

The multijet background (MJ) due to the heavy flavor quark decays and pions faking electrons is hard to be accurately simulated by MC
-> Derive from the data-driven estimation method.

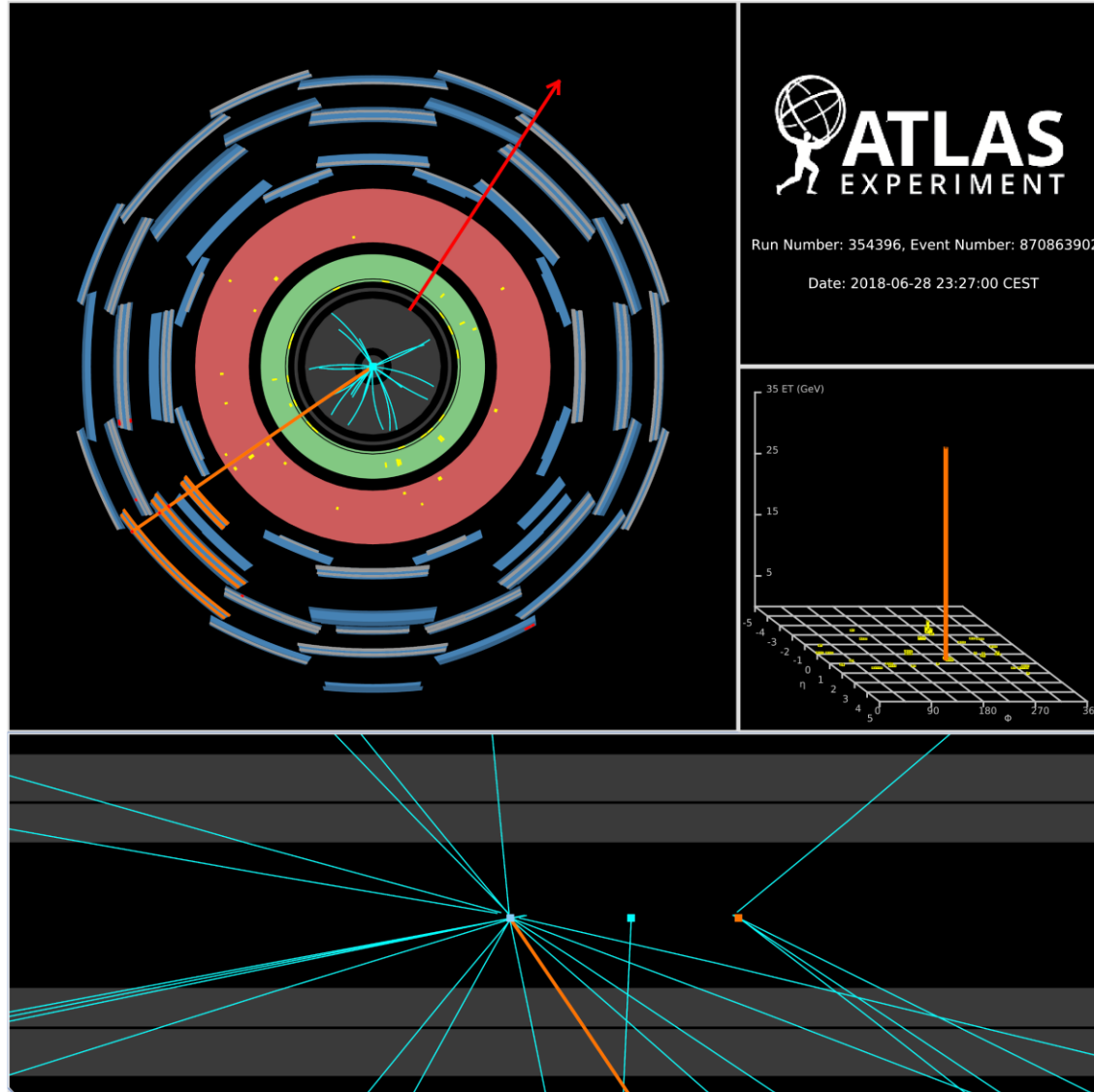
Two elements to be estimated:

- MJ yield in the signal region.
- MJ shape in the signal region.

Unfolding



Event display of the low pile-up data



The event display of a W^- boson candidate at 13 TeV.

- Orange line: muon
- Red arrow: missing transverse momentum

Event kinematics:

- $p_T^\mu = 35$ GeV
- Reconstructed $p_T^W = 16$ GeV
- $m_T = 77$ GeV
- $p_T^{\text{miss}} = 49$ GeV.