



# Dispelling the $\sqrt{L}$ myth for the High-Luminosity LHC

Alberto Belvedere

ETTORE MAJORANA FOUNDATION AND CENTRE FOR SCIENTIFIC CULTURE  
INTERNATIONAL SCHOOL OF SUBNUCLEAR PHYSICS  
Erice - 17th June 2024

# Introduction

During phase-I LHC will collect roughly  $300 \text{ fb}^{-1}$  while at the end HL-LHC  $3000 \text{ fb}^{-1}$  should be available, **a factor 10 increase in statistics.**

It is generally believed that the sensitivity  $S$  on a parameter of interest scales as  $\sqrt{L}$  because:

$$S \simeq \frac{S}{\sqrt{B}} \simeq \sqrt{L} \frac{\sigma_S}{\sigma_B}$$

The  $\sqrt{L}$  **scaling** is overly **conservative at best**, and **unrealistic in practice.**

This can be proven in at least four different examples:

- Measurement of **4 top quarks production cross-section.**
- Measurement of the **top mass** in hadronic decays of boosted top quarks.
- Constrain **BSM** parameters **in electroweak processes** involving the **top quark.**
- Constrain **BSM** parameters **in Higgs boson** physics.

# Introduction

It is generally believed that the sensitivity  $S$  on a parameter of interest scales as  $\sqrt{L}$  because:

$$S \simeq \frac{S}{\sqrt{L}}$$

**Focus of this talk.**

(Other examples explained  
in [arXiv:2402.07985](https://arxiv.org/abs/2402.07985))

$$\delta \approx \sqrt{S} \approx \sqrt{L}$$

During phase-I LHC will be available, a factor 10 increase in statistics. At the end HL-LHC 3000  $fb^{-1}$  should be available, a factor 10 increase in statistics.

The  $\sqrt{L}$  scaling is overly conservative at best, and unrealistic in practice.

This can be proven in at least four different examples:

- Measurement of **4 top quarks production cross-section**.
- Measurement of the **top mass** in in hadronic decays of boosted top quarks.
- Constrain **BSM** parameters **in electroweak processes** involving the **top quark**.
- Constrain **BSM** parameters **in Higgs boson** physics.

# The Standard Model Effective Field Theory (SMEFT)

It is possible to Taylor expand the SM Lagrangian in terms of  $\frac{E}{\Lambda}$ :

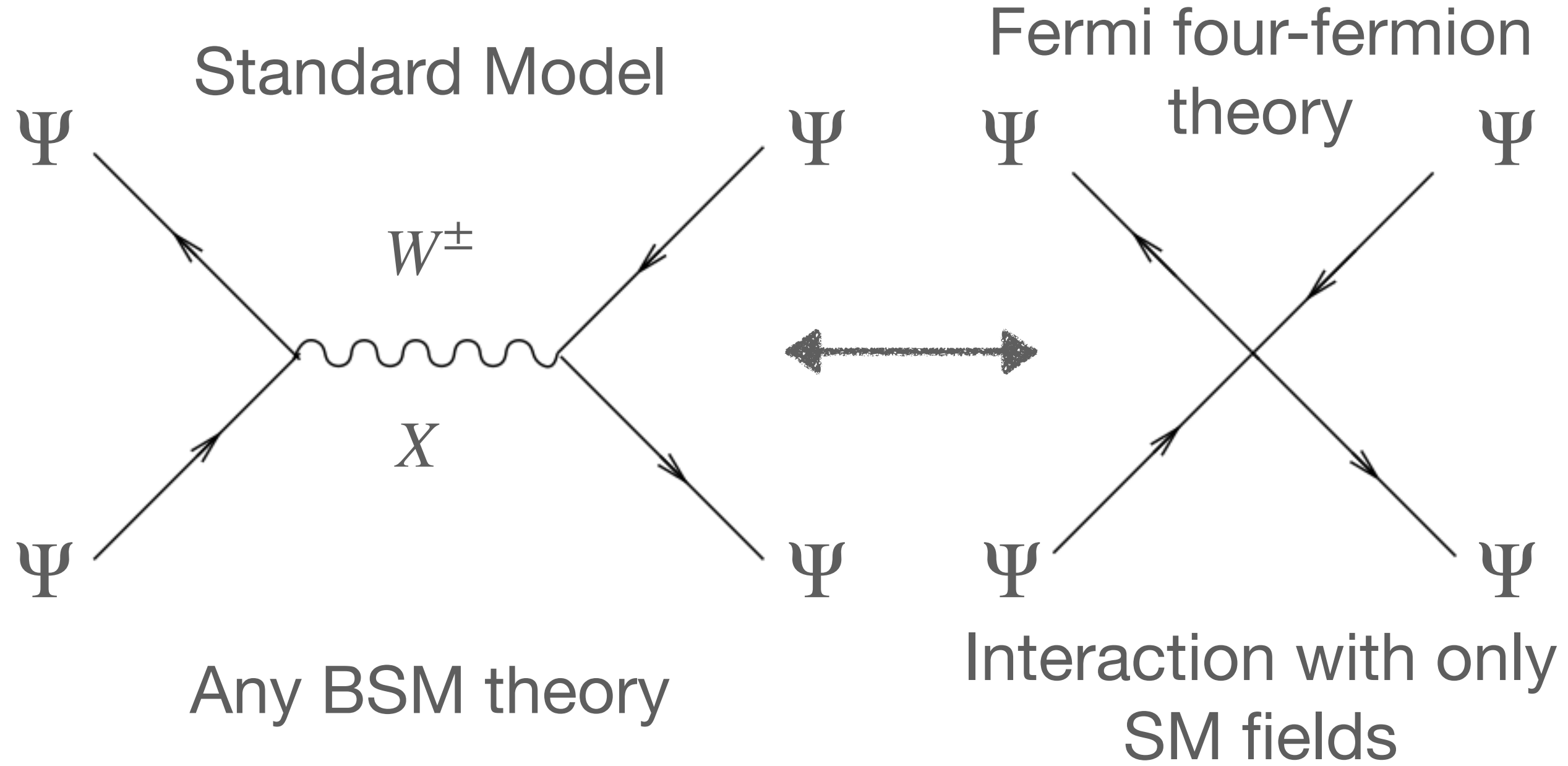
$$\mathcal{L} = \frac{1}{\Lambda^0} \mathcal{L}_{SM}^{(d=4)} + \frac{1}{\Lambda} \mathcal{L}^{(d=5)} + \frac{1}{\Lambda^2} \mathcal{L}^{(d=6)} + \dots$$

$$\mathcal{L}^{(d)} = \sum_i c_i \mathcal{O}_i^{(d)}$$

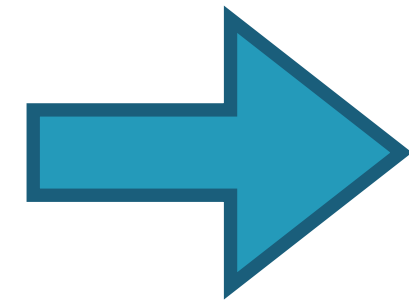
All the possible local interactions allowed by symmetries.

- $c_i$  are called Wilson coefficients - **free parameters** of the model.
- $\mathcal{O}_i^{(d)}$  represent all the possible **operators**.

# SMEFT: a different approach to new physics



- The SM can be considered as an effective low-energy theory.
- BSM physics can be thought of as modifications of the interactions containing only SM fields.



Completely **model independent**

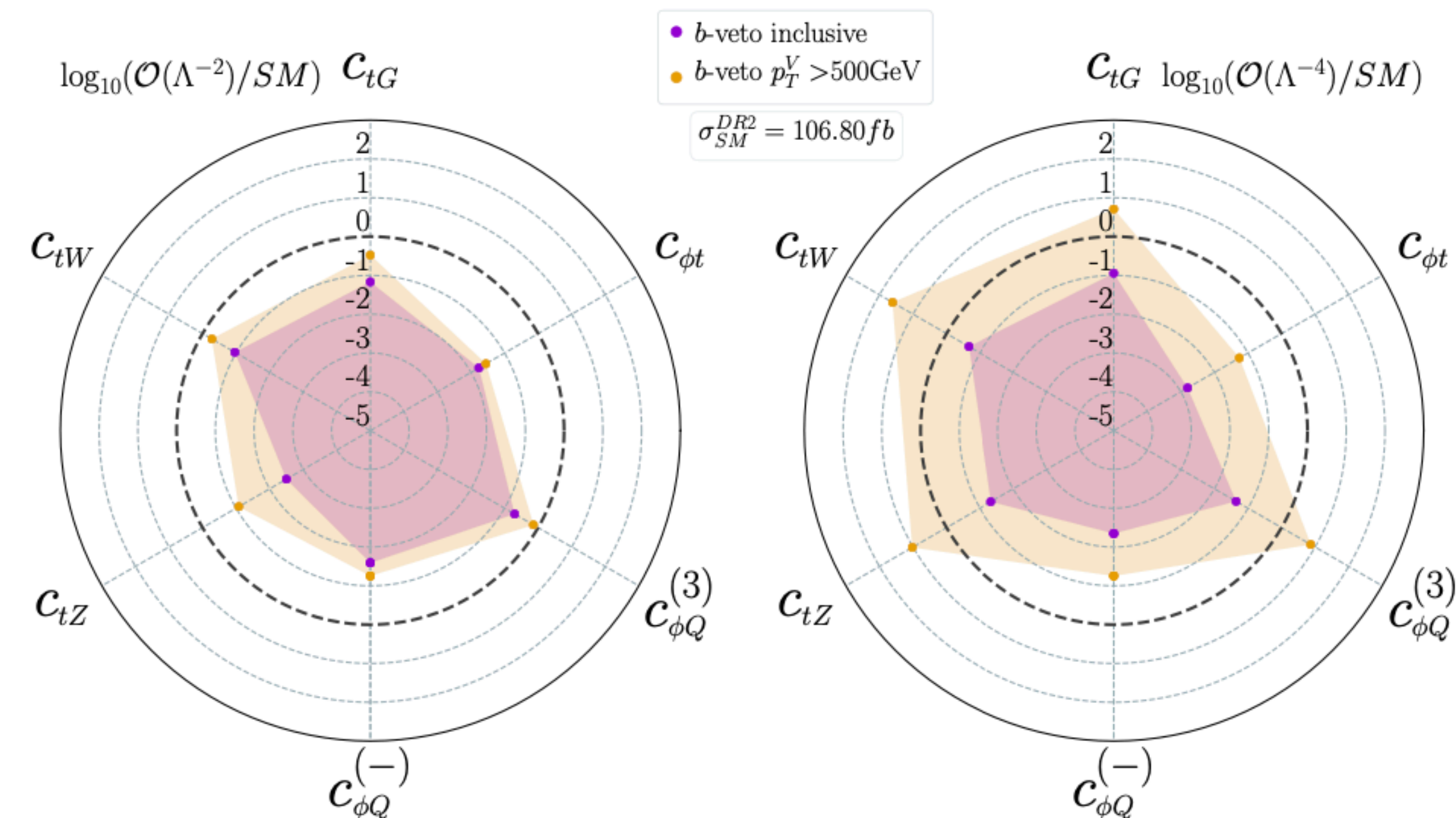
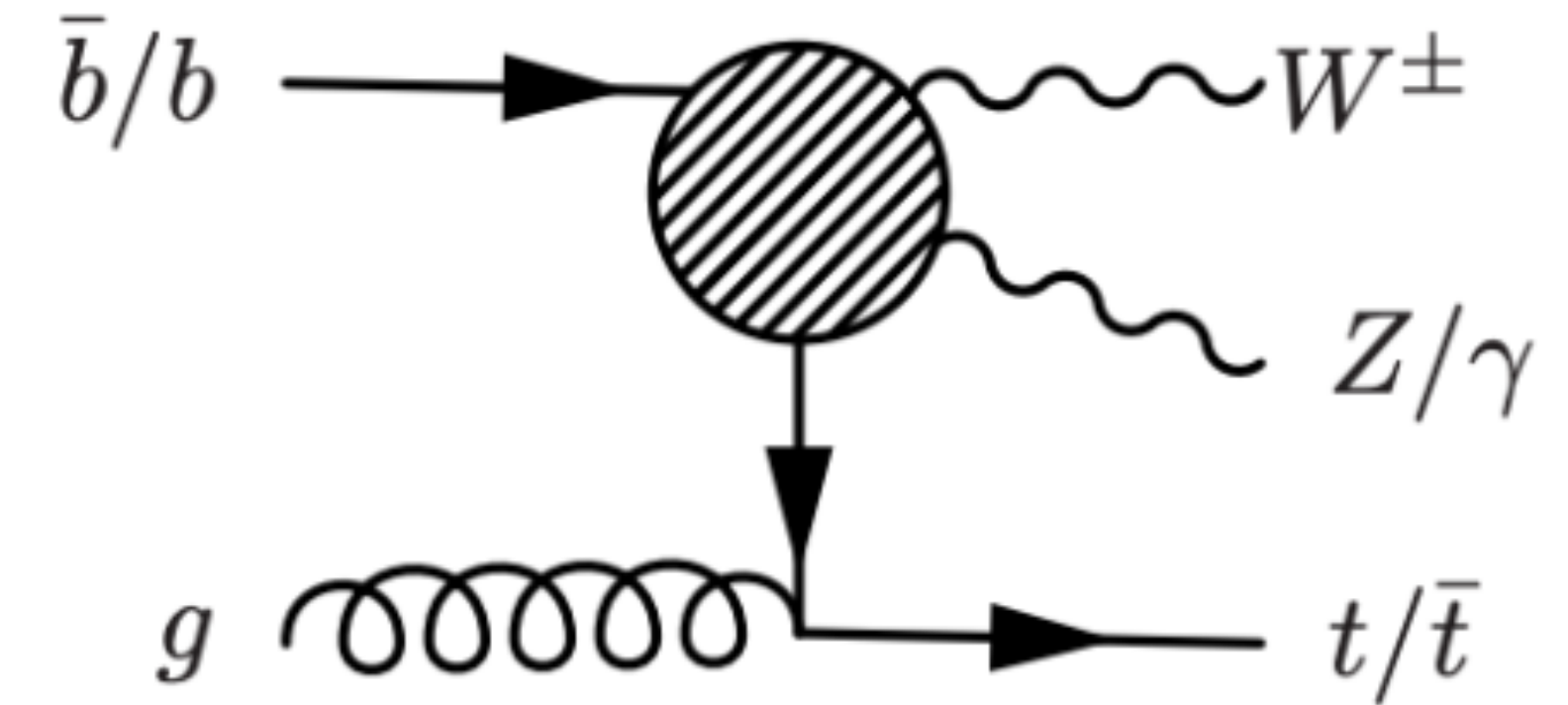


# The tWZ process

- tWZ has **never been observed**, however CMS has recently [published](#) the first evidence.
- Very challenging process:
  - ▶ **Small** predicted **cross-section** ( $\sigma_{tWZ} = 136 \text{ fb}$ ).
  - ▶ **Large** irreducible **background** ( $t\bar{t}Z$  process):
- Its discovery would allow the analysis of SMEFT operators only loosely constrained with possible new physics arising from **bW**→**tZ** vertices.
- This would lead in **energy growth** for some specific operators.

[arXiv:1904.05637](#)


[arXiv:2111.03080](#)



# Simulation of EFT predictions

$$\mathcal{L} = \frac{1}{\Lambda^0} \mathcal{L}_{SM}^{(d=4)} + \frac{1}{\Lambda} \mathcal{L}^{(d=5)} + \frac{1}{\Lambda^2} \mathcal{L}^{(d=6)} + \dots$$

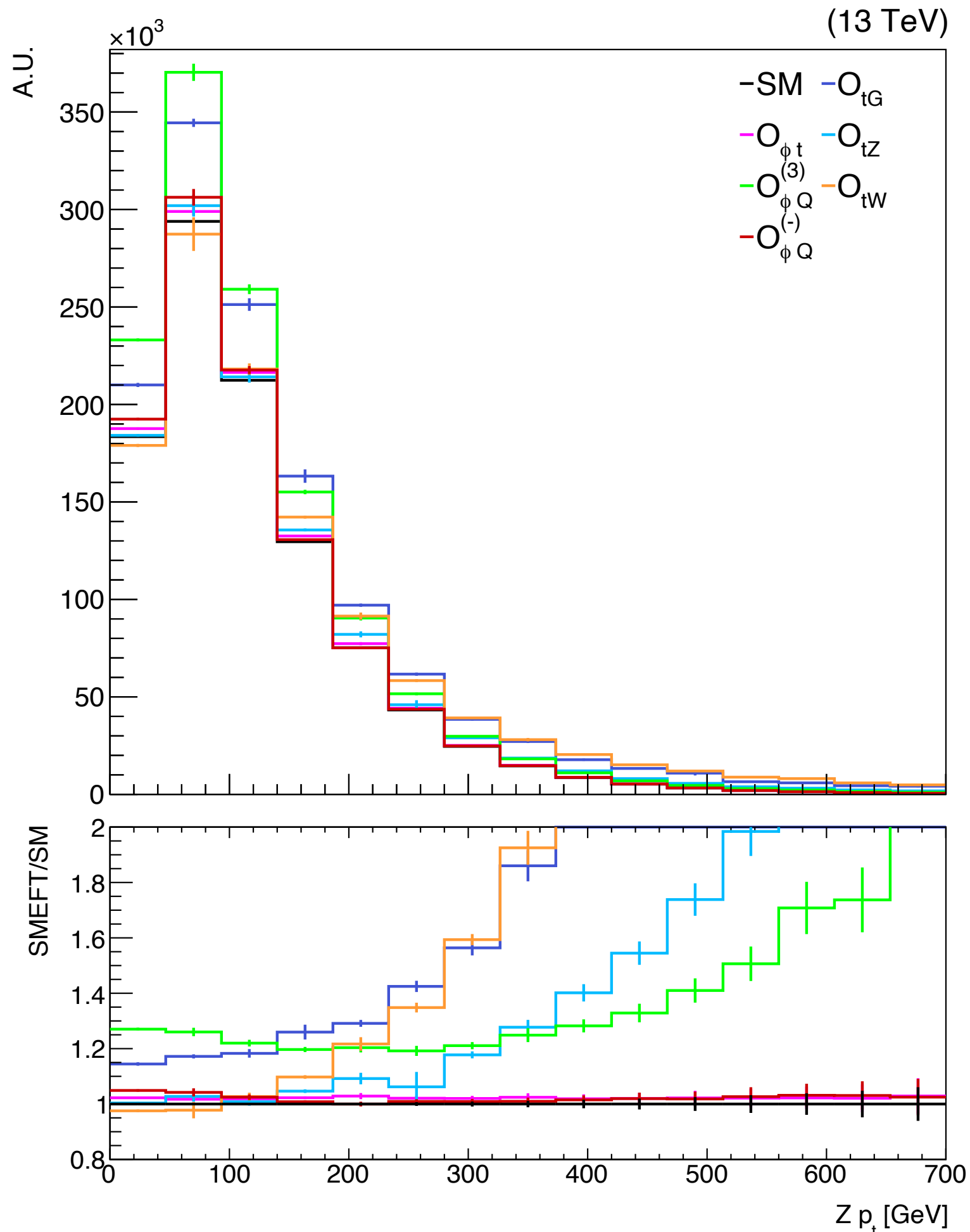
- MC samples generated with MadGraph5 using the **SMEFT@NLO** model.
- Production of samples to test the **effect of the operators** on relevant variables of the process.
- Event showered with **PYTHIA8**.
- Operators are tested setting as value for the corresponding Wilson coefficients one.
- $\Lambda$  is kept fixed at the value of 1 TeV.


$$\sum_i \frac{c_i \mathcal{O}_i^{(6)}}{\Lambda^2}$$

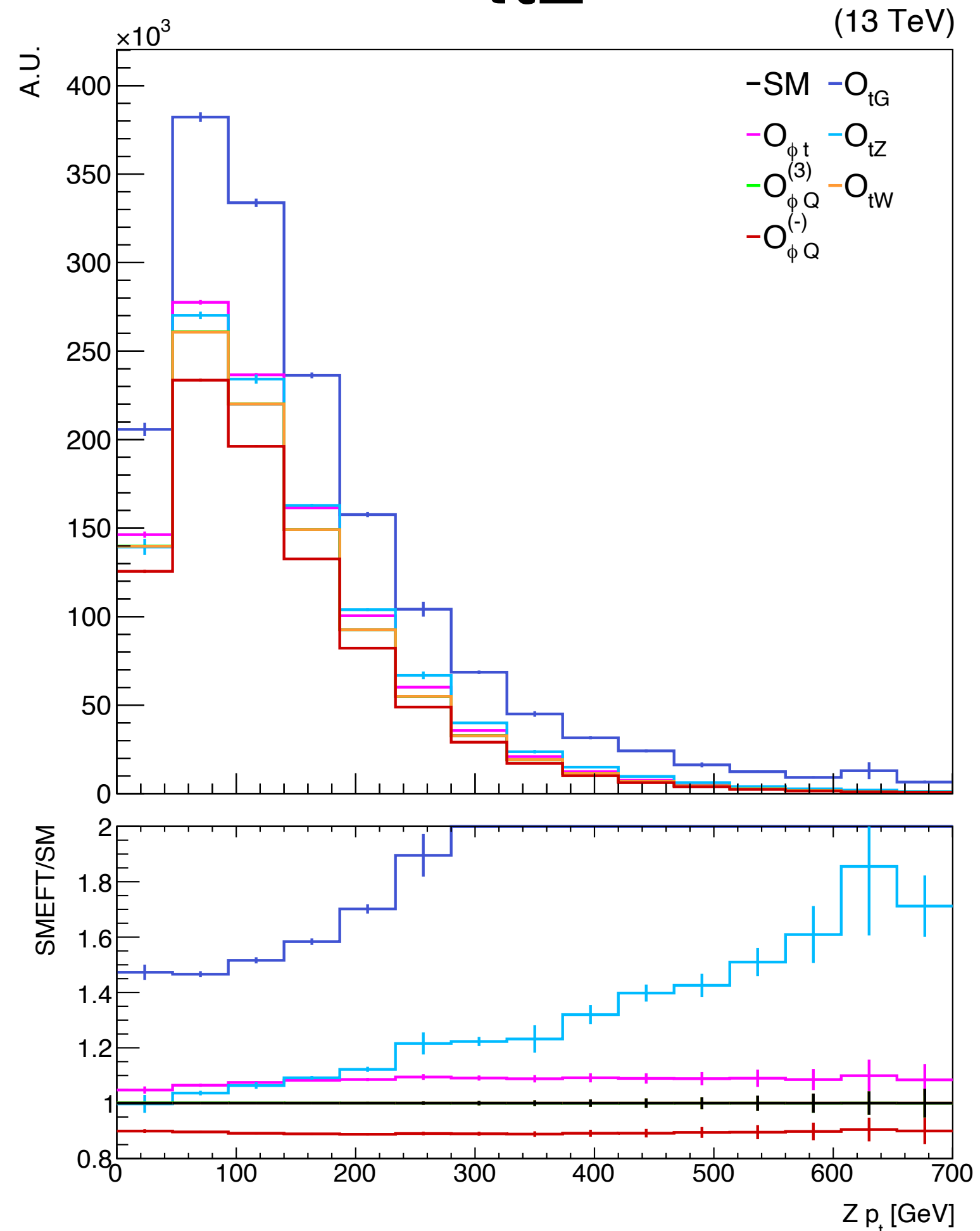
# $Z p_t$ distributions

Comparison between SM distributions and SMEFT distributions obtained through reweighting.

**tWZ**



**t $\bar{t}$ Z**



- Excess for the  $\mathcal{O}_{\phi Q}^{(3)}$  operator shows energy growth in **tWZ**.
- Also the  $\mathcal{O}_{tZ}$  operator shows energy growth for both **tWZ** and **t $\bar{t}$ Z**.
- **tWZ** is more sensitive than **t $\bar{t}$ Z** to most of the operators.



# The CMS measurement and EFT interpretation

The **first evidence** of the **tWZ process** was recently [reported](#) by the CMS experiment.

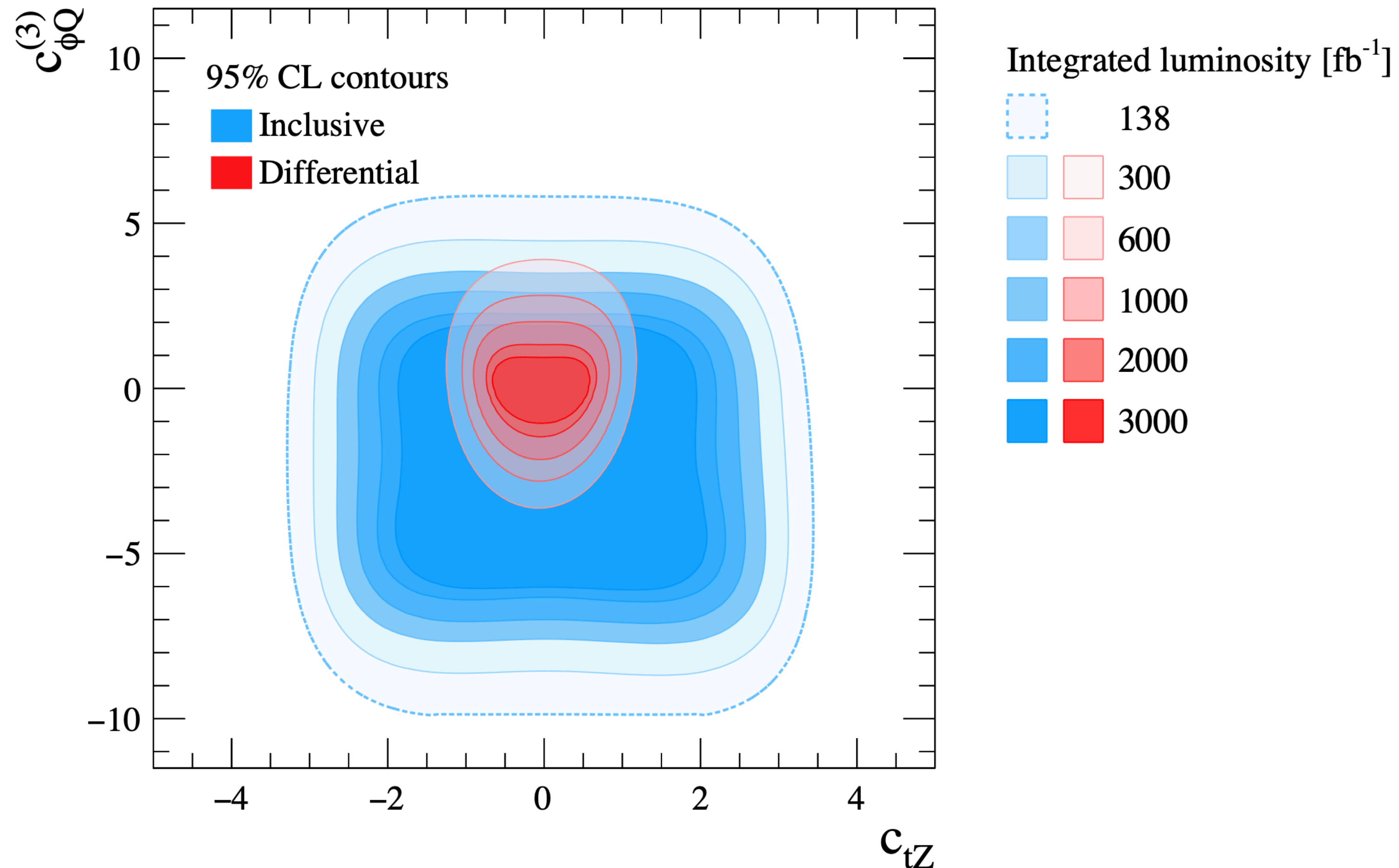
Analysis performed on  $138 \text{ fb}^{-1}$  in 3 and 4 leptons final states using ML algorithms to distinguish the **tWZ** process from the background, mostly **t $\bar{t}$ Z**.

Observed (expected) **significance** =  $3.5\sigma$  ( $1.4\sigma$ )

$$\sigma_{tWZ} = 0.37 \pm 0.05(\text{stat}) \pm 0.10(\text{syst}) \text{ pb}$$

Building on this result, we were able to compute limits on the Wilson coefficients of the  $\mathcal{O}_{\phi Q}^{(3)}$  and  $\mathcal{O}_{tZ}$  operators.

# EFT limits and perspectives for HL-LHC



Collecting enough data to perform a **differential measurement in  $\mathbf{Z} p_T$**  would allow to set **much tighter constraints** on the Wilson coefficients.

# Conclusion

When projecting the sensitivity to a parameter of interest in scenarios with larger integrated luminosity, **the sensitivity is often scaled by  $\sqrt{L}$** .

However the **sensitivity improvement** from higher statistics is not only related to the increase in the number of events:

- Improved estimation and constraint of **systematics uncertainties**.
- **More advanced analysis strategies**.

With the HL-LHC, we will be able to **challenge the Standard Model with unprecedented precision**, offering the potential for possible surprises.