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# Flavor physics

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# General remarks

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- I have to make assumptions about what you know
- Please ask questions (in class and outside)
- Email: yg73@cornell.edu
- Much more in
  - Book: <https://www.classe.cornell.edu/~yuvalg/>
  - TASI lectures: "Just a taste" 1711.03624
  - Online lectures on the SM and on Flavor
- The plan:
  - Flavor in the SM
  - Flavor beyond the SM

# Some data

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$$\text{Br}(K_L \rightarrow \mu^+ \mu^-) [s \rightarrow d \mu^+ \mu^-] = 6.84 \times 10^{-9}$$

$$\text{Br}(K^- \rightarrow \mu^- \bar{\nu}) [s \rightarrow u \mu^- \bar{\nu}] = 0.6356$$

$$\text{Br}(D^+ \rightarrow \bar{K}^0 e^+ \nu) [c \rightarrow s e^+ \nu] = 8.82 \times 10^{-2}$$

$$\text{Br}(D^+ \rightarrow \bar{K}^0 \mu^+ \nu) [c \rightarrow s \mu^+ \nu] = 8.74 \times 10^{-2}$$

$$\text{Br}(B \rightarrow X_c e \nu) [b \rightarrow c e \nu] = 0.1086$$

$$\text{Br}(B_s \rightarrow \mu^+ \mu^-) [b \rightarrow s \mu^+ \mu^-] = 2.4 \times 10^{-9}$$

$$\text{Br}(B^- \rightarrow D^0 \mu^- \bar{\nu}) [b \rightarrow c \mu^- \bar{\nu}] = 2.27 \times 10^{-2}$$

$$\text{Br}(B^- \rightarrow \pi^0 \mu^- \bar{\nu}) [b \rightarrow u \mu^+ \bar{\nu}] = 7.80 \times 10^{-5}$$

● What patterns do you see?

# What we learn from the data

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- **Lepton universality.** Swapping one generation of leptons with another does not appear to affect the branching ratios of these transitions.
  - **Flavor-changing neutral currents are small.** Charge-neutral transitions are suppressed compared to transitions between hadrons of different charge.
  - **Generation hierarchy.** Decays between third and first generation are suppressed compared to that of third to second generation.
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The hope is to understand these within the SM

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# What is HEP?

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Find the basic laws of Nature

More formally

$$\mathcal{L} = ?$$

- We have quite a good answer
- It is very elegant, it is based on axioms and symmetries
- We use particles to answer this question

# Building Lagrangians

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- We built a Lagrangian by the following input
  - The symmetries we impose
  - The fields and their transformations
- The output is the Lagrangian such that
  - It is the most general that obeys the symmetries
  - We truncate it at some order, usually fourth

# Example: A SM

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(i) The symmetry is a local

$$SU(3)_C \times SU(2)_L \times U(1)_Y$$

(ii) There are three fermion generations (QUDLE)

$$Q_L(3, 2)_{+1/6}, \quad U_R(3, 1)_{2/3}, \quad D_R(3, 1)_{-1/3},$$

$$L_L(1, 2)_{-1/2}, \quad E_R(1, 1)_{-1}$$

(iii) There is a single scalar multiplet

$$\phi(1, 2)_{+1/2}$$



# $\mathcal{L}$ for a SM

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The most general  $\mathcal{L}$  is given by

$$\mathcal{L} = \mathcal{L}_{\text{kin}} + \mathcal{L}_{\psi} + \mathcal{L}_{\phi} + \mathcal{L}_{\text{Yuk}}$$

- No mass term for the fermions:  $\mathcal{L}_{\psi} = 0$
- Kinetic terms (with gauge interactions)
- The scalar potential that lead to SSB

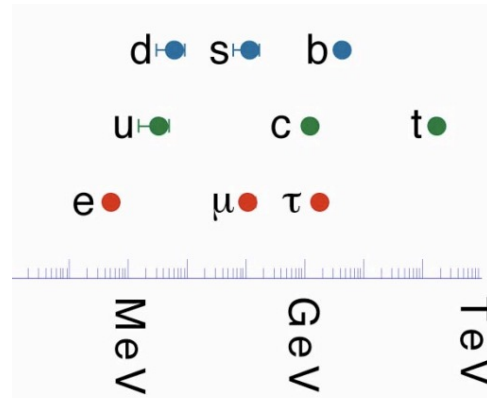
$$-\mathcal{L}_{\phi} = -\mu^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2$$

- The Yukawas that lead to fermion masses

$$\mathcal{L}_{\text{Yuk}} \sim Y^U Q U \phi + Y^D Q D \phi + Y^{\ell} L E \phi$$

# A SM vs The SM

- “A SM” is the theory without the values of the parameters
- “The SM” is the one we have with a given set of values for the parameters
- The hierarchy of masses is a "The SM" thing



- "A SM" and even more "The SM" are very delicate. The fact that they work is far from trivial

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# The SM: Flavor physics

# Moving between the bases

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Flavor is about moving between the bases



## Vocabulary

- Flavor basis: The couplings to the  $W$  are diagonal
- Mass basis: Where the mass matrix is diagonal

# Basis choice

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Consider the following three matrices

- Down-type mass
- Up-type mass
- Coupling to the  $W$

We can at most diagonalize two out of the three matrices

- We usually choose the coupling to the  $W$  to be the non-diagonal one
- For the leptons  $m_\nu$  vanishes so we have only two matrices to deal with

# The CKM matrix

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In the mass basis we have

$$\mathcal{L}_W \sim V_{ij} u_i d_j W$$

- $V$  is called the CKM matrix
- The point is that we cannot have  $m_U$ ,  $m_D$  and the couplings to the  $W$  to be diagonal in the same basis
- In the mass basis the  $W$  interaction change flavor, that is, flavor is not conserved

# The CKM matrix

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The CKM matrix is

- Unitary
- Four parameters: three angles and a phase

$$|V_{us}|, \quad |V_{cb}|, \quad |V_{ub}|, \quad \delta_{KM}$$

$\delta_{KM}$  is related to CPV

- Close to a unit matrix

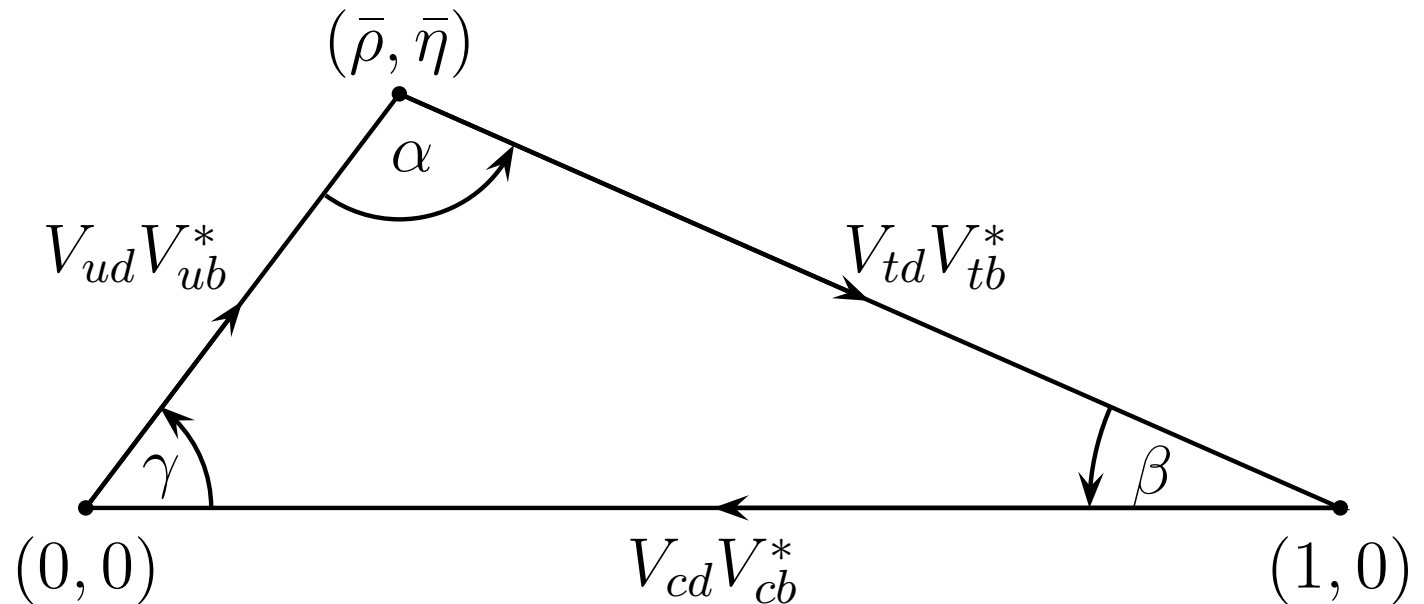
$$|V| \approx \begin{pmatrix} 0.97383 & 0.2272 & 3.96 \times 10^{-3} \\ 0.2271 & 0.97296 & 4.221 \times 10^{-2} \\ 8.14 \times 10^{-3} & 4.161 \times 10^{-2} & 0.99910 \end{pmatrix}$$

# The Unitarity Triangle

- Since the CKM is unitary its elements satisfy

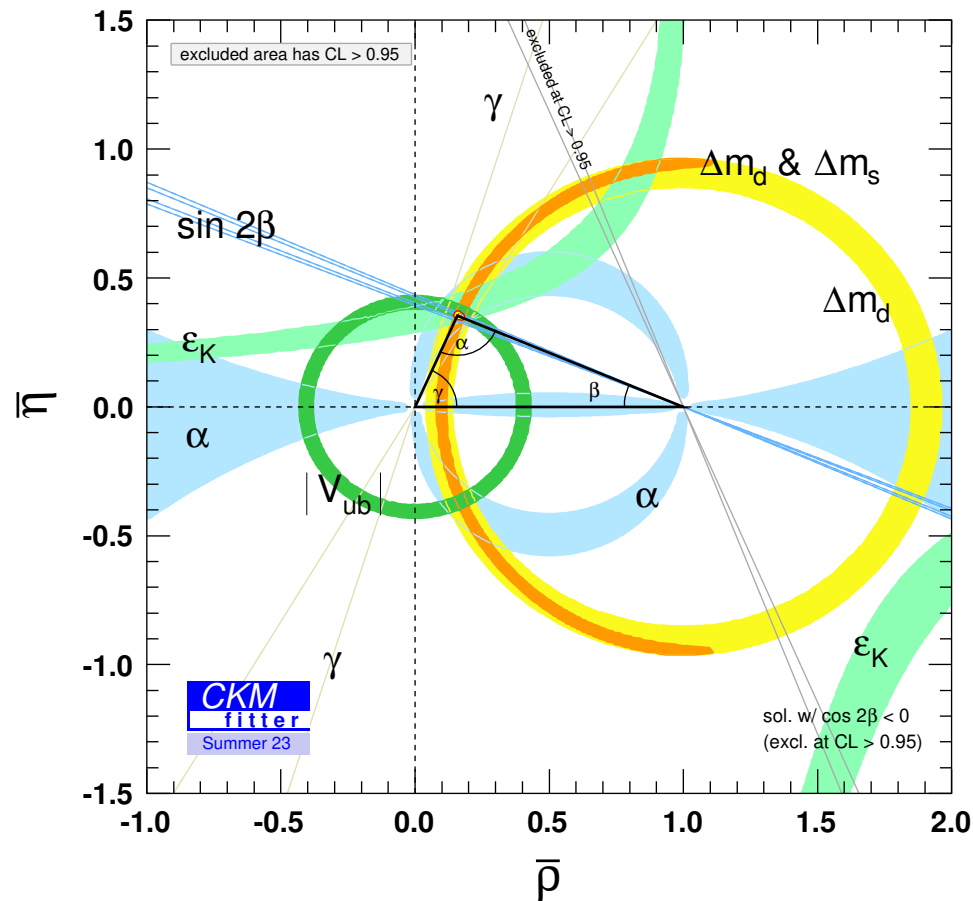
$$\sum_k V_{ik} V_{jk}^* = \delta_{ij}$$

- A nice way to describe it





# The SM works



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# FCNC

# FCNC

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Very important concept in flavor physics

- FCNC=Flavor Changing Neutral Current
  - FCCC=Flavor Changing Charged Current
  - For example,  $b \rightarrow cl\nu$  vs  $b \rightarrow sl^+\ell^-$
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- Vocabulary: diagonal couplings vs universal couplings

# FCNC in Nature

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- In Nature FCNC are highly suppressed
  - Historically,  $K \rightarrow \mu\nu$  vs  $K_L \rightarrow \mu\mu$
  - The suppression was also seen in  $c$  and  $b$  decays
- In the SM there are no FCNC at tree level. Very nice!
  - In the SM we have four neutral bosons,  $g, \gamma, Z, h$ . Their couplings are diagonal
  - The reasons why they are diagonal, and what it takes to have FCNC, is not always trivial
  - Of course we have FCNC at one loop (two charged current interactions give a neutral one)

# Photon and gluon tree level FCNC

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- For exact gauge interactions the couplings are always diagonal. It is part of the kinetic term

$$\partial_\mu \delta_{ij} \rightarrow (\partial_\mu + iqA_\mu) \delta_{ij}$$

- Symmetries are nice...
- In any extension of the SM the photon and gluon couplings are flavor diagonal

# Higgs tree level FCNC

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- The Higgs is a possible source of FCNC
- With one Higgs doublet, the mass matrix is aligned with the Yukawa.  $\phi \sim v + H$

$$\mathcal{L}_m \sim Y_{ij} v \bar{d}_L^i d_R^j \quad \mathcal{L}_{int} \sim Y_{ij} H \bar{d}_L^i d_R^j$$

- With two doublets we could have tree level FCNC

# $Z$ exchange FCNC

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- For broken gauge symmetry there is no FCNC when:  
“All the fields with the same QN in the unbroken symmetry also have the same QN in the broken part”
- In the SM the  $Z$  coupling is diagonal since all  $q = -1/3$   
RH quarks are  $(3, 1)_{-1/3}$  under  $SU(2) \times U(1)$
- Adding quarks with different representations can generate tree level FCNC  $Z$  couplings, like  $b'_L(3, 1)_{-1/3}$
- Same condition for new neutral gauge bosons (usually denoted by  $Z'$ )

# Tree level FCNC in “a SM”

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- In a generic model we expect tree-level FCNC
- In a SM there are specific reasons for not having it
  - All fermions with the same charge also have the same hypercharge
  - There is only one Higgs doublet
- We need to be careful when extending the SM



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# FCNC at one loop

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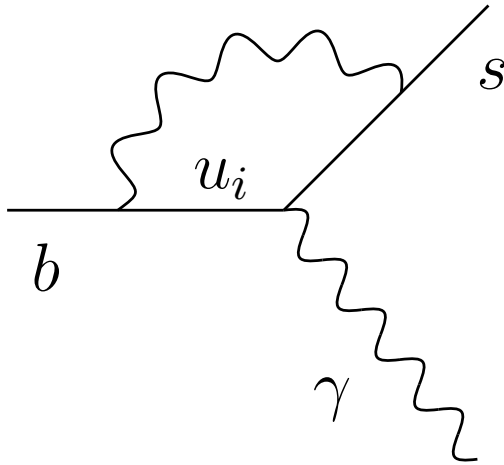
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- We understand why FCNC are suppressed in “a SM”:  
There is no tree level exchange
- Yet, there are more suppression factors in “The SM”
  - CKM factors
  - Mass ratio factors: GIM mechanism
- The loop factor in “a SM” is universal
- The other factors in “the SM” are not universal, they depend on the quarks that are involved

# Loop: example

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$$A(b \rightarrow s\gamma) \propto \sum V_{ib} V_{is}^*$$



What is  $\sum V_{ib} V_{is}^*$ ?

# GIM Mechanism

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What we really have is

$$A(b \rightarrow s\gamma) \propto \sum V_{ib} V_{is}^* f(m_i)$$

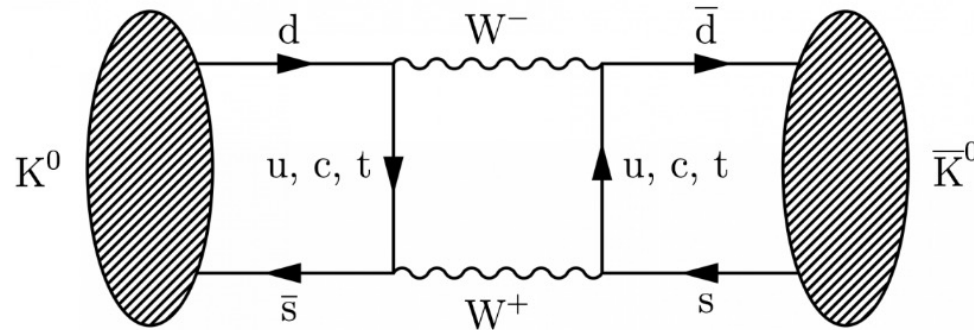
- Unitarity of the CKM  $\Rightarrow m_i$  independent term in  $f$  vanishes  $\Rightarrow$  The amplitude depend on the masses
- For  $m_i \ll m_W^2$  we have

$$A \sim \frac{m_i^2}{m_W^2}$$

- In  $s$  decays it gives  $m_c^2/m_W^2$  extra suppression
- In charm it gives  $m_s^2/m_W^2$ extra suppression
- Numerically, not important for  $b$  decays,  $m_t \sim m_W$

# Example: kaon mixing

Roughly it is giving by box diagram

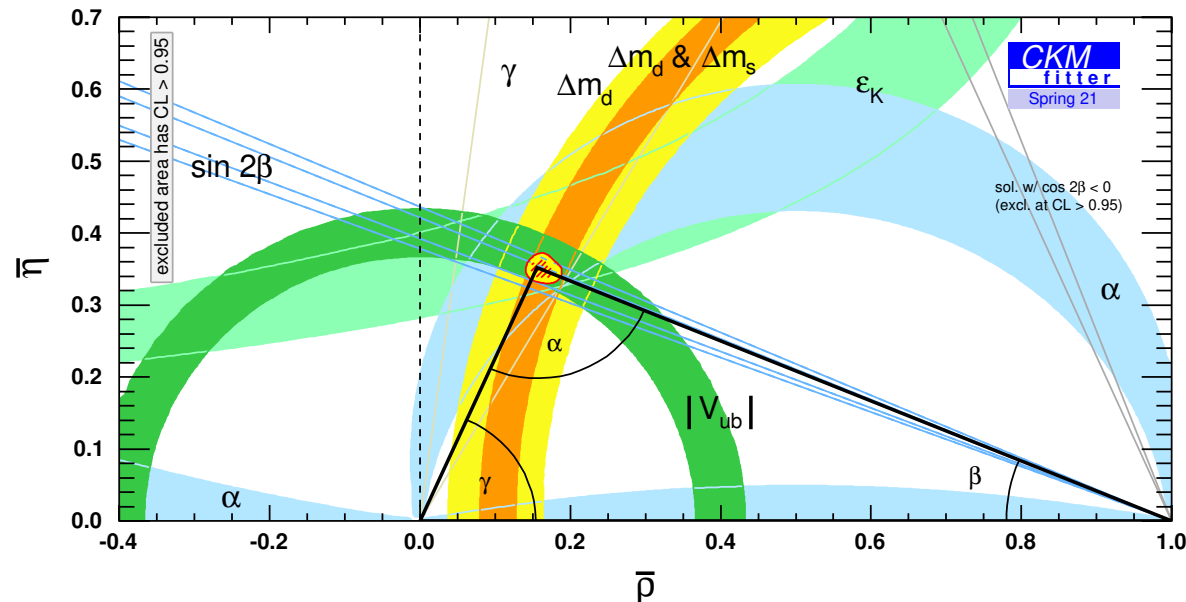


● The result is

$$\Delta m_K \sim \frac{g^4}{16\pi^2} \times \frac{m_c^2}{m_W^2} \times |V_{us}V_{ud}|^2 \sim 10^{-8}$$

● Different factors for different mesons

# Summary: The SM is very special



- “A SM” is special
- “The SM” is even more special

The fact that the data confirm the SM is far from trivial